

From Eq.(2.25), we can obtain the Fourier transform of  $\psi$  as

$$\psi = \begin{cases} f^+(k_x, k_y) e^{k_y y} & (y \geq 0) \\ f^-(k_x, k_y) e^{k_y y} & (y \leq 0) \end{cases} \quad (2.26)$$

Then, in order to satisfy the radiation condition at  $y \rightarrow \pm\infty$ , we must choose

$$k_y^- = -k_y^+ = \begin{cases} j\sqrt{k_x^2 + k_y^2 - k^2} & (k \leq \sqrt{k_x^2 + k_y^2}) \\ -\sqrt{k^2 - k_x^2 - k_y^2} & (k \geq \sqrt{k_x^2 + k_y^2}) \end{cases} \quad (2.27)$$

As in the above aperture radiation example, we make the Fourier transform of the field components in Eq.(2.19) after setting  $A_{z0} = 0$  and  $A_x = \psi$  and obtain

$$\hat{H}_y^+ = \frac{j k_x^2}{\mu_0} \psi^+, \quad \hat{E}_x^+ = -\frac{k_x k_y}{j \omega \mu_0 \epsilon_0} \psi^+, \quad (2.28a)$$

$$\hat{H}_y^- = -\frac{j k_x^2}{\mu_0} \psi^-, \quad \hat{E}_x^- = -\frac{k_x^2 k_y}{j \omega \mu_0 \epsilon_0} \psi^-, \quad (2.28b)$$

$$\hat{H}_z^+ = 0, \quad \hat{E}_z^+ = \frac{1}{j \omega \mu_0 \epsilon_0} [k^2 - k_x^2] \psi^+. \quad (2.28c)$$

According to the continuity of the tangential electric fields, i.e.,

$E_x^+|_{y=0} = E_x^-|_{y=0}$  and  $E_z^+|_{y=0} = E_z^-|_{y=0}$ , we have

$$f^+(k_x, k_y) = f^-(k_x, k_y) = f(k_x, k_y). \quad (2.29)$$

Also, according to the boundary conditions of the magnetic fields at  $y=0$ , i.e.,  $H_x^+ - H_x^- = -J_x$  or  $\hat{H}_y^+ - \hat{H}_y^- = -\hat{J}_y$  when  $y=0$ , by using Eq.(2.26) and considering Eqs.(2.25) and (2.27),  $f(k_x, k_y)$  can be determined by

$$f(k_x, k_y) = \frac{j \mu_0}{2 k_y^+} \hat{J}_y, \quad (2.30)$$

# Advanced Electromagnetic Theory

**Hyo J. Eom**



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